STAR FORMATION

Hauyu Baobab Liu (呂浩宇) Department of Physics, National Sun Yat-sen University

STAR FORMATION

Hauyu Baobab Liu (呂浩宇) Department of Physics, National Sun Yat-sen University



STATUS OF THIS FIELD

You can be known for discovering or addressing something fundamental

- 2. There are more questions than theories and answers
- 3. The duty-cycle of making your own hypothesis and then test it is short
 - 4. This field is under-credited and is presently not particularly hot \otimes

Why Star-Formation Matters



Why Star-Formation Matters



Stars Make Galaxies and the Parent Dark Halo Visible to us



(1) Halo Occupation Distribution (2) Kennicutt-Schmidt Law (3) Gao-Solomon Relation

(1) <u>Halo Occupation Distribution</u> (2) Kennicutt-Schmidt Law (3) Gao-Solomon Relation

- (i) How sub-halos (candidates of galaxies ?) accrete baryonic material? And how feedback and other physics quench the accretion (e.g., AGN feedback, ram pressure stripping in a galaxy cluster). The observations of inter-galactic medium (IGM) are to address these issues.
- (ii) How the first-generation stars (also known as the Population III stars) formed, and what were their roles in shaping the visible universe?



(1) Halo Occupation Distribution (2) Kennic

(2) <u>Kennicutt-Schmidt Law</u> (3) Gao-Solomon Relation



(1) Halo Occupation Distribution (2) Kennicutt-Sch

(2) Kennicutt-Schmidt Law (3) Gao-Solomon Relation



https://en.wikipedia.org/wiki/Initial mass function

(1) Halo Occupation Distribution (2) <u>Kennicutt-Schmidt Law</u> (3) Gao-Solomon Relation



Kuiper et al. 1938, ApJ, 88, 472

(1) Halo Occupation Distribution (2) Kennicutt

(2) Kennicutt-Schmidt Law (3) Gao-Solomon Relation



$$\Sigma_{SFR} = 2.5 \times 10^{-4} \left(\frac{\Sigma_{gas}}{1 \, M_{\odot} \, pc^{-2}} \right)^{1.4} M_{\odot} yr^{-1} kpc^{-2}$$

Kennicutt et al. 1998, ApJ, 498, 541

Explain 1: free-fall timescale

Star-formation volume density

$$\rho_{SFR} \propto \frac{\rho_{gas}}{t_{ff}}$$

$$\propto \frac{\rho_{gas}}{\left(G\rho_{gas}\right)^{-0.5}} \propto \rho_{gas}^{1.5}$$

Assuming constant scale-height

$$\Sigma_{SFR} \propto
ho_{SFR} \Sigma_{gas} \propto
ho_{gas}$$

(1) Halo Occupation Distribution (2) Kennicutt-

(2) Kennicutt-Schmidt Law (3) Gao-Solomon Relation



$$\Sigma_{SFR} = 2.5 \times 10^{-4} \left(\frac{\Sigma_{gas}}{1 \, M_{\odot} \, pc^{-2}} \right)^{1.4} M_{\odot} yr^{-1} kpc^{-2}$$

Kennicutt et al. 1998, ApJ, 498, 541

Explain 2: galactic dynamical timescale

$$\Sigma_{SFR} \propto \frac{\Sigma_{gas}}{\tau_{dyn}} \propto \Sigma_{gas} \Omega_{gas}$$

 Ω_{gas} : local orbital timescale

(1) Halo Occupation Distribution

(2) <u>Kennicutt-Schmidt Law</u> (3) Gao-Solomon Relation



(1) Halo Occupation Distribution (2) Kennicutt-Schmidt Law (3) Gao-Solomon Relation



Gao & Solomon 2004, ApJ, 606, 271

$$SFR = (1.8 \times 10^{-8}) \times \left(\frac{M_{dense}}{1 M_{\odot}}\right) M_{\odot} yr^{-1}$$

(1) Halo Occupation Distribution (2) Kennicutt-

(2) Kennicutt-Schmidt Law (3) G

(3) Gao-Solomon Relation





$$SFR = (1.8 \times 10^{-8}) \times \left(\frac{M_{dense}}{1 M_{\odot}}\right) M_{\odot} yr^{-1}$$

Again this law implies very low star-forming efficiency. In addition, the definition of dense gas remains ambiguous (Jiao et al. submitted).



(1) Halo Occupation Distribution (2) Kennicu

(2) Kennicutt-Schmidt Law

(3) Gao-Solomon Relation



Gao & Solomon 2004, ApJ, 606, 271

$$SFR = (1.8 \times 10^{-8}) \times \left(\frac{M_{dense}}{1 M_{\odot}}\right) M_{\odot} yr^{-1}$$

This Law is very strange. Ten molecular clouds of $10^5 M_{\odot}$ of gas mass and one molecular cloud with $10^6 M_{\odot}$ of gas mass consume the same amount of gas mass to starformation every year, in spite that the stars they form are very different (Jiao, Xu, Liu et al. in prep.)



Stars Form in Clusters (Multiplicity is Essential)



Stellar Initial Mass Function





https://en.wikipedia.org/wiki/Initial_mass_function

 $M_{max} - M_{ecl}$ relation

Weidner, Kroupa, Bonnell 2010, MNRAS, 401, 275







A Simplified Picture of Interstellar Medium

```
1 pc = 3 \times 10^{16} meters
```







A Simplified Picture of Interstellar Medium



A Simplified Picture of Interstellar Medium



Equation of Continuity:

<u>mass:</u>

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

momentum:

$$\frac{\partial(\rho v_i)}{\partial t}\hat{\imath} + \vec{\nabla} \cdot \left((\rho v_i)\vec{v}\right)\hat{\imath} = 0$$



Momentum density $p_i = \rho v_i$

Equation of Continuity:

<u>mass:</u>

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

$$\frac{\partial(\rho v_i)}{\partial t}\hat{\imath} + \vec{\nabla} \cdot \left((\rho v_i)\vec{v}\right)\hat{\imath} = 0 \qquad \text{Momentum density } p_i = \rho v_i$$

$$\frac{\partial(\rho)}{\partial t}v_i + \rho \frac{\partial(v_i)}{\partial t} + (\partial_j \rho)v_i v_j + \rho(\partial_j v_i)v_j + \rho v_i(\partial_j v_j) = 0$$



Equation of Continuity:

<u>mass:</u>

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

$$\frac{\partial(\rho v_i)}{\partial t}\hat{\imath} + \vec{\nabla} \cdot \left((\rho v_i)\vec{v}\right)\hat{\imath} = 0 \qquad \text{Momentum density } p_i = \rho v_i$$

$$\frac{\partial(\rho)}{\partial t}v_{i} + \rho \frac{\partial(v_{i})}{\partial t} + (\partial_{j}\rho)v_{i}v_{j} + \rho(\partial_{j}v_{i})v_{j} + \rho v_{i}(\partial_{j}v_{j}) = 0$$

$$\vec{\nabla}\rho$$



Equation of Continuity:

<u>mass:</u>

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

$$\frac{\partial(\rho v_i)}{\partial t}\hat{\imath} + \vec{\nabla} \cdot \left((\rho v_i)\vec{v}\right)\hat{\imath} = 0 \qquad \text{Momentum density } p_i = \rho v_i$$

$$\frac{\partial(\rho)}{\partial t}v_{i} + \rho \frac{\partial(v_{i})}{\partial t} + (\partial_{j}\rho)v_{i}v_{j} + \rho(\partial_{j}v_{i})v_{j} + \rho v_{i}(\partial_{j}v_{j}) = 0$$

$$\vec{\nabla}\rho$$

$$v_{i}(\vec{\nabla}\rho \cdot \vec{v})$$



Equation of Continuity:

<u>mass:</u>

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

$$\frac{\partial(\rho v_i)}{\partial t}\hat{\imath} + \vec{\nabla} \cdot \left((\rho v_i)\vec{v}\right)\hat{\imath} = 0 \qquad \text{Momentum density } p_i = \rho v_i$$

$$\frac{\partial(\rho)}{\partial t}v_{i} + \rho \frac{\partial(v_{i})}{\partial t} + (\partial_{j}\rho)v_{i}v_{j} + \rho(\partial_{j}v_{i})v_{j} + \rho v_{i}(\partial_{j}v_{j}) = 0$$

$$\vec{\nabla}\rho \qquad \rho(\vec{v} \cdot (\vec{\nabla}v_{i})) \qquad \rho v_{i}(\vec{\nabla} \cdot \vec{v})$$

$$v_{i}(\vec{\nabla}\rho \cdot \vec{v})$$



Equation of Continuity:

<u>mass:</u>

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

$$\frac{\partial(\rho v_i)}{\partial t}\hat{\imath} + \vec{\nabla} \cdot ((\rho v_i)\vec{v})\hat{\imath} = 0 \qquad \text{Momentum density } p_i = \rho v_i$$

$$\frac{\partial(\rho)}{\partial t}\boldsymbol{v}_{i} + \rho \frac{\partial(\boldsymbol{v}_{i})}{\partial t} + \boldsymbol{v}_{i} (\vec{\nabla}\rho \cdot \vec{v}) + \rho (\vec{v} \cdot (\vec{\nabla}\boldsymbol{v}_{i})) + \rho \boldsymbol{v}_{i} (\vec{\nabla} \cdot \vec{v}) = 0$$



Equation of Continuity:

<u>mass:</u>

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

$$\frac{\partial(\rho v_i)}{\partial t}\hat{\imath} + \vec{\nabla} \cdot \left((\rho v_i)\vec{v}\right)\hat{\imath} = 0 \qquad \text{Momentum density } p_i = \rho v_i$$

$$\frac{\partial(\rho)}{\partial t}\boldsymbol{v}_{i} + \rho \frac{\partial(\boldsymbol{v}_{i})}{\partial t} + \boldsymbol{v}_{i} (\vec{\nabla}\rho \cdot \vec{v}) + \rho (\vec{v} \cdot (\vec{\nabla}\boldsymbol{v}_{i})) + \rho \boldsymbol{v}_{i} (\vec{\nabla} \cdot \vec{v}) = 0$$
$$\boldsymbol{v}_{i} \left[\frac{\partial(\rho)}{\partial t} + (\vec{\nabla}\rho \cdot \vec{v}) + \rho (\vec{\nabla} \cdot \vec{v}) \right] + \rho \left[\frac{\partial(\boldsymbol{v}_{i})}{\partial t} + \vec{v} \cdot (\vec{\nabla}\boldsymbol{v}_{i}) \right] = 0$$



Equation of Continuity:

<u>mass:</u>

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

$$\underline{\text{momentum:}}$$

$$\frac{\partial (\rho v_i)}{\partial t} \hat{\imath} + \vec{\nabla} \cdot ((\rho v_i) \vec{v}) \hat{\imath} = 0 \qquad \text{Momentum density } p_i = \rho v_i$$

$$\frac{\partial (\rho)}{\partial t} v_i + \rho \frac{\partial (v_i)}{\partial t} + v_i (\vec{\nabla} \rho \cdot \vec{v}) + \rho (\vec{v} \cdot (\vec{\nabla} v_i)) + \rho v_i (\vec{\nabla} \cdot \vec{v}) = 0$$

$$v_i \left[\frac{\partial (\rho)}{\partial t} + (\vec{\nabla} \rho \cdot \vec{v}) + \rho (\vec{\nabla} \cdot \vec{v}) \right] + \rho \left[\frac{\partial (v_i)}{\partial t} + \vec{v} \cdot (\vec{\nabla} v_i) \right] = 0$$





Equation of Continuity (force free):

<u>mass:</u>

$$----\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

$$\frac{\partial(\rho v_i)}{\partial t}\hat{\imath} + \vec{\nabla} \cdot \left((\rho v_i)\vec{v}\right)\hat{\imath} = 0 \qquad \text{Momentum density } p_i = \rho v_i$$
$$\frac{\partial(\rho)}{\partial t}v_i + \rho \frac{\partial(v_i)}{\partial t} + v_i(\vec{\nabla}\rho \cdot \vec{v}) + \rho(\vec{v} \cdot (\vec{\nabla}v_i)) + \rho v_i(\vec{\nabla} \cdot \vec{v}) = 0$$
$$v_i \left[\frac{\partial(\rho)}{\partial t} + (\vec{\nabla}\rho \cdot \vec{v}) + \rho(\vec{\nabla} \cdot \vec{v})\right] + \rho \left[\frac{\partial(v_i)}{\partial t} + \vec{v} \cdot (\vec{\nabla}v_i)\right] = 0$$

Equation of Continuity (force free):

<u>mass:</u>

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

$$\rho\left[\frac{\partial(\boldsymbol{v}_i)}{\partial t} + \vec{v} \cdot (\vec{\nabla}\boldsymbol{v}_i)\right] = 0$$



Equation of Continuity (with pressure *P* and gravitational acceleration \vec{g}): <u>mass</u>:

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

$$\frac{\partial(\boldsymbol{\rho}\boldsymbol{v}_{i})}{\partial t}\hat{\imath} + \vec{\nabla}\cdot\left((\boldsymbol{\rho}\boldsymbol{v}_{i})\vec{v}\right)\hat{\imath} = \left[-\left(\vec{\nabla}P\right)_{i} + \boldsymbol{\rho}g_{i}\right]\hat{\imath}$$



Equation of Continuity (with pressure *P* and gravitational acceleration \vec{g}): <u>mass</u>:

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

momentum:

$$\frac{\partial(\rho v_i)}{\partial t}\hat{\imath} + \vec{\nabla} \cdot \left((\rho v_i)\vec{v}\right)\hat{\imath} = \left[-\left(\vec{\nabla}P\right)_i - \rho\left(\vec{\nabla}\phi\right)_i\right]\hat{\imath}$$

Poisson Equation $abla^2 \phi = 4\pi G \rho$



Equation of Continuity (with pressure *P* and gravitational acceleration \vec{g}): <u>mass</u>:

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

momentum:

$$\boldsymbol{\rho}\left[\frac{\partial(\boldsymbol{v}_i)}{\partial t} + \vec{v} \cdot (\vec{\nabla}\boldsymbol{v}_i)\right] \hat{\boldsymbol{\iota}} = \left[-\left(\vec{\nabla}\boldsymbol{P}\right)_i - \boldsymbol{\rho}\left(\vec{\nabla}\boldsymbol{\phi}\right)_i\right] \hat{\boldsymbol{\iota}}$$

Poisson Equation $abla^2 \phi = 4\pi G \rho$



Equation of Continuity (with pressure *P* and gravitational acceleration \vec{g}): <u>mass</u>:

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

momentum:

$$\frac{\partial(\boldsymbol{v}_i)}{\partial t} + \vec{v} \cdot (\vec{\nabla} \boldsymbol{v}_i) = -\frac{\left(\vec{\nabla} P\right)_i}{\rho} - \left(\vec{\nabla} \phi\right)_i$$

Poisson Equation $abla^2 \phi = 4\pi G \rho$



Equation of Continuity (with pressure *P* and gravitational acceleration \vec{g}): <u>mass</u>:

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

momentum:

$$\frac{\partial(\boldsymbol{v}_{i})}{\partial t} + \vec{v} \cdot (\vec{\nabla}\boldsymbol{v}_{i}) = -\frac{\left(\vec{\nabla}P\right)_{i}}{\rho} - \left(\vec{\nabla}\phi\right)_{i}$$

Perturbation theory

$$\begin{bmatrix} \rho = \rho_0 + \rho_1, & \rho_0 \gg \rho_1, & \vec{\nabla}\rho_0 = 0 \text{ [uniform initial condition]} \\ P = P_0 + P_1, & P_0 \gg P_1, & \vec{\nabla}P_0 = 0 \text{ [uniform initial condition]} \\ \vec{v} = \vec{v}_0 + \vec{v}_1, & \vec{v}_0 = 0 \text{ [initially quiescent cloud]} \\ \phi = \phi_0 + \phi_1, & \phi_0 \gg \phi_1 \end{bmatrix}$$

Plugging into the two equations of continuity, and the Poisson equation

Poisson Equation $abla^2 \phi = 4\pi G \rho$

Equation of state $P = c_s^2 \rho$, $c_s \equiv isothermal sound speed$

$$\begin{bmatrix} \frac{\partial \rho_1}{\partial t} + \rho_0 \vec{\nabla} \cdot \vec{v_1} = 0 \\ \rho_0 \frac{\partial \vec{v_1}}{\partial t} = -c_s^2 \vec{\nabla} \rho_1 - \rho_0 \vec{\nabla} \phi_1 \\ \nabla^2 \phi_1 = 4\pi G \rho_1 \end{bmatrix}$$



$$\begin{bmatrix} \frac{\partial \rho_1}{\partial t} + \rho_0 \vec{\nabla} \cdot \vec{v_1} = 0 & (1) \\ \rho_0 \frac{\partial \vec{v_1}}{\partial t} = -c_s^2 \vec{\nabla} \rho_1 - \rho_0 \vec{\nabla} \phi_1 & (2) \\ \nabla^2 \phi_1 = 4\pi G \rho_1 & (3) \end{bmatrix}$$



$$\frac{\partial}{\partial t}(1) \Longrightarrow \frac{\partial^2 \rho_1}{\partial t^2} + \rho_0 \vec{\nabla} \cdot \frac{\partial \vec{v_1}}{\partial t} = 0$$

$$\begin{bmatrix} \frac{\partial \rho_1}{\partial t} + \rho_0 \vec{\nabla} \cdot \vec{v_1} = 0 & (1) \\ \rho_0 \frac{\partial \vec{v_1}}{\partial t} = -c_s^2 \vec{\nabla} \rho_1 - \rho_0 \vec{\nabla} \phi_1 & (2) \\ \nabla^2 \phi_1 = 4\pi G \rho_1 & (3) \\ \frac{\partial}{\partial t} (1) \Rightarrow \frac{\partial^2 \rho_1}{\partial t^2} + \rho_0 \vec{\nabla} \cdot \frac{\partial \vec{v_1}}{\partial t} = 0 \\ \Rightarrow \frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 - \rho_0 \nabla^2 \phi_1 = 0 & (4) \\ \Rightarrow \frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 - 4\pi G \rho_0 \rho_1 = 0 & (5) \\ \end{bmatrix}$$

Dispersion relation $\frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 = 4\pi G \rho_0 \rho_1$



$$\begin{bmatrix} \frac{\partial \rho_1}{\partial t} + \rho_0 \vec{\nabla} \cdot \vec{v_1} = 0 & (1) \\ \rho_0 \frac{\partial \vec{v_1}}{\partial t} = -c_s^2 \vec{\nabla} \rho_1 - \rho_0 \vec{\nabla} \phi_1 & (2) \\ \nabla^2 \phi_1 = 4\pi G \rho_1 & (3) \\ \frac{\partial}{\partial t} (1) \Rightarrow \frac{\partial^2 \rho_1}{\partial t^2} + \rho_0 \vec{\nabla} \cdot \frac{\partial \vec{v_1}}{\partial t} = 0 \\ \Rightarrow \frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 - \rho_0 \nabla^2 \phi_1 = 0 & (4) \\ \Rightarrow \frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 - 4\pi G \rho_0 \rho_1 = 0 & (5) \\ \end{bmatrix}$$



Dispersion relation $\frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 = 4\pi G \rho_0 \rho_1$

let $\rho_1 \propto e^{i(\vec{k}\cdot\vec{x}-\omega t)}$

$$\Rightarrow -\omega^2 + k^2 c_s^2 = 4\pi G \rho_0$$
$$\Rightarrow \omega^2 = k^2 c_s^2 - 4\pi G \rho_0$$

$$\begin{bmatrix} \frac{\partial \rho_1}{\partial t} + \rho_0 \vec{\nabla} \cdot \vec{v_1} = 0 & (1) \\ \rho_0 \frac{\partial \vec{v_1}}{\partial t} = -c_s^2 \vec{\nabla} \rho_1 - \rho_0 \vec{\nabla} \phi_1 & (2) \\ \nabla^2 \phi_1 = 4\pi G \rho_1 & (3) \\ \frac{\partial}{\partial t} (1) \Rightarrow \frac{\partial^2 \rho_1}{\partial t^2} + \rho_0 \vec{\nabla} \cdot \frac{\partial \vec{v_1}}{\partial t} = 0 \\ \Rightarrow \frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 - \rho_0 \nabla^2 \phi_1 = 0 & (4) \\ \Rightarrow \frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 - 4\pi G \rho_0 \rho_1 = 0 & (5) \\ \end{bmatrix}$$



Dispersion relation $\frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 = 4\pi G \rho_0 \rho_1$

let
$$\rho_1 \propto e^{i(\vec{k}\cdot\vec{x}-\omega t)}$$

$$\Rightarrow -\omega^2 + k^2 c_s^2 = 4\pi G \rho_0$$
$$\Rightarrow \omega^2 = k^2 c_s^2 - 4\pi G \rho_0$$

define
$$k_J = \sqrt{\frac{4\pi G \rho_0}{c_S^2}}$$

$$\frac{\partial \rho_{1}}{\partial t} + \rho_{0} \vec{\nabla} \cdot \vec{v_{1}} = 0 \qquad (1)$$

$$\rho_{0} \frac{\partial \vec{v_{1}}}{\partial t} = -c_{s}^{2} \vec{\nabla} \rho_{1} - \rho_{0} \vec{\nabla} \phi_{1} \qquad (2) \qquad (2) \qquad (3)$$

$$\vec{\nabla}^{2} \phi_{1} = 4\pi G \rho_{1} \qquad (3) \qquad (4) \qquad (3) \qquad (4) \qquad (3) \qquad (4) \qquad (3) \qquad (4) \qquad (4) \qquad (4) \qquad (4) \qquad (5) \qquad (4) \qquad (5) \qquad (4) \qquad (5) \qquad (5)$$



Dispersion relation $\frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 = 4\pi G \rho_0 \rho_1$

let
$$\rho_1 \propto e^{i(\vec{k}\cdot\vec{x}-\omega t)}$$

$$\Rightarrow -\omega^2 + k^2 c_s^2 = 4\pi G \rho_0$$
$$\Rightarrow \omega^2 = k^2 c_s^2 - 4\pi G \rho_0$$

define
$$k_J = \sqrt{\frac{4\pi G \rho_0}{c_s^2}}$$

When $\begin{bmatrix} k < k_J, & \omega^2 < 0, & \text{perturbation grows exponentially} \\ k > k_J, & \omega^2 > 0, & \text{perturbation oscillates} \end{bmatrix}$

define $k_J = \sqrt{\frac{4\pi G \rho_0}{c_s^2}}$ (Jeans length)



When $\begin{cases} k < k_J, & \omega^2 < 0, & \text{perturbation grows exponentially} \\ k > k_J, & \omega^2 > 0, & \text{perturbation oscillates} \end{cases}$

define $k_J = \sqrt{\frac{4\pi G \rho_0}{c_s^2}}$ (Jeans length)



When $\begin{cases} k < k_J, & \omega^2 < 0, & \text{perturbation grows exponentially} \\ k > k_J, & \omega^2 > 0, & \text{perturbation oscillates} \end{cases}$

$$k = \frac{2\pi}{\lambda}$$

define
$$k_J = \sqrt{\frac{4\pi G \rho_0}{c_s^2}}$$
 (Jeans length)

When $\begin{bmatrix} k < k_J, & \omega^2 < 0, & \text{perturbation grows exponentially} \\ k > k_J, & \omega^2 > 0, & \text{perturbation oscillates} \end{bmatrix}$

 $k = \frac{2\pi}{\lambda}$ \implies perturbation grows exponentially when $\lambda > \lambda_J \equiv \frac{2\pi}{k_J} = c_s \sqrt{\frac{\pi}{G\rho_0}}$ (Jeans length)



define
$$k_J = \sqrt{\frac{4\pi G \rho_0}{c_s^2}}$$
 (Jeans length)

When $\begin{cases} k < k_J, & \omega^2 < 0, & \text{perturbation grows exponentially} \\ k > k_J, & \omega^2 > 0, & \text{perturbation oscillates} \end{cases}$

$$k = \frac{2\pi}{\lambda} \implies \text{perturbation grows exponentially when}$$
$$\lambda > \lambda_J \equiv \frac{2\pi}{k_J} = c_s \sqrt{\frac{\pi}{G\rho_0}} \text{ (Jeans length)}$$
$$\implies \text{mass of perturbation}$$
$$M_J \equiv \frac{4}{3} \pi \left(\frac{\lambda_J}{2}\right)^3. \text{ (Jeans mass)}$$

An initially uniform and quiescent molecular cloud will fragment into substructures of Jeans masses. The spatial separations of these sub-structures are approximately the Jeans length.



define
$$k_J = \sqrt{\frac{4\pi G \rho_0}{c_s^2}}$$
 (Jeans length)

When $\begin{cases} k < k_J, & \omega^2 < 0, & \text{perturbation grows exponentially} \\ k > k_J, & \omega^2 > 0, & \text{perturbation oscillates} \end{cases}$

$$k = \frac{2\pi}{\lambda} \implies \text{perturbation grows exponentially when}$$
$$\lambda > \lambda_J \equiv \frac{2\pi}{k_J} = c_s \sqrt{\frac{\pi}{G\rho_0}} \text{ (Jeans length)}$$
$$\implies \text{mass of perturbation}$$
$$M_J \equiv \frac{4}{3} \pi \left(\frac{\lambda_J}{2}\right)^3 \rho_0 \qquad \text{(Jeans mass)}$$

An initially uniform and quiescent molecular cloud will fragment into substructures of Jeans masses. The spatial separations of these sub-structures are approximately the Jeans length.

An initially uniform and quiescent molecular cloud will just collapse to form a black hole.



(1) Linearization works OK with the 1st order equations, but does not make sense at the 0th order

$$\begin{bmatrix} \rho = \rho_0 + \rho_1, & \rho_0 \gg \rho_1, & \vec{\nabla}\rho_0 = 0 \text{ [uniform initial condition]} \\ P = P_0 + P_1, & P_0 \gg P_1, & \vec{\nabla}P_0 = 0 \text{ [uniform initial condition]} \\ \vec{v} = \vec{v}_0 + \vec{v}_1, & \vec{v}_0 = 0 \text{ [initially quiescent cloud]} \\ \phi = \phi_0 + \phi_1, & \phi_0 \gg \phi_1 \end{cases}$$

$$\frac{\partial(\boldsymbol{v}_i)}{\partial t} + \vec{v} \cdot (\vec{\nabla} \boldsymbol{v}_i) = -\frac{\left(\vec{\nabla} P\right)_i}{\rho} - \left(\vec{\nabla} \phi\right)_i \qquad \text{Oth order momentum eq.}$$

(1) Linearization works OK with the 1st order equations, but does not make sense at the 0th order

$$\begin{bmatrix} \rho = \rho_0 + \rho_1, & \rho_0 \gg \rho_1, & \vec{\nabla}\rho_0 = 0 \text{ [uniform initial condition]} \\ P = P_0 + P_1, & P_0 \gg P_1, & \vec{\nabla}P_0 = 0 \text{ [uniform initial condition]} \\ \vec{v} = \vec{v}_0 + \vec{v}_1, & \vec{v}_0 = 0 \text{ [initially quiescent cloud]} \\ \phi = \phi_0 + \phi_1, & \phi_0 \gg \phi_1 \end{bmatrix}$$

$$\frac{\partial(\boldsymbol{v}_i)}{\partial t} + \vec{v} \cdot (\vec{\nabla} \boldsymbol{v}_i) = -\frac{\left(\vec{\nabla} P\right)_i}{\rho} - \left(\vec{\nabla} \phi\right)_i \qquad \text{Oth order momentum eq.}$$

$$\Rightarrow \left(\vec{\nabla} P_0 \right)_i = 0 = \rho_0 \left(\vec{\nabla} \phi_0 \right)_i$$
$$\Rightarrow \vec{\nabla} \phi_0 = 0 \Rightarrow \nabla^2 \phi = 0 \neq 4\pi G \rho_0$$

An initially quiescent cloud cannot exist. A finite-sized cloud will undergo gravitational contraction.

(1) Linearization works OK with the 1st order equations, but does not make sense at the 0th order

$$\begin{split} \rho &= \rho_0 + \rho_1, \quad \rho_0 \gg \rho_1, \quad \vec{\nabla} \rho_0 = 0 \quad \text{[uniform initial condition]} \\ P &= P_0 + P_1, \quad P_0 \gg P_1, \quad \vec{\nabla} P_0 = 0 \quad \text{[uniform initial condition]} \\ \vec{v} &= \vec{v}_0 + \vec{v}_1, \quad \vec{v}_0 = 0 \quad \text{[initially quiescent cloud]} \\ \phi &= \phi_0 + \phi_1, \quad \phi_0 \gg \phi_1 \end{split}$$

$$\frac{\partial(\boldsymbol{v}_i)}{\partial t} + \vec{v} \cdot (\vec{\nabla} \boldsymbol{v}_i) = -\frac{\left(\vec{\nabla} P\right)_i}{\rho} - \left(\vec{\nabla} \phi\right)_i \qquad \text{Oth order momentum eq.}$$

$$\Rightarrow \left(\vec{\nabla} P_0 \right)_i = 0 = \rho_0 \left(\vec{\nabla} \phi_0 \right)_i$$
$$\Rightarrow \vec{\nabla} \phi_0 = 0 \Rightarrow \nabla^2 \phi = 0 \neq 4\pi G \rho_0$$

An initially quiescent cloud cannot exist. A finite-sized cloud will undergo gravitational contraction.

(2) The growing timescale is minimized when $|\omega^2|$ is maximized at k = 0

Global collapse has a shorter characteristic $\omega^2 = k^2 c_s^2 - 4\pi G \rho_0$ timescale than perturbation growth

An initially uniform molecular cloud cannot fragment to form a cluster of star

An initially uniform molecular cloud cannot fragment to form a cluster of star

To make cloud fragmentation efficient, we need to make the global free-fall timescale larger than the local free-fall timescale

$$t_{ff} = \sqrt{\frac{3\pi}{16G\rho}}$$

An initially uniform molecular cloud cannot fragment to form a cluster of star

To make cloud fragmentation efficient, we need to make the global free-fall timescale larger than the local free-fall timescale

$$t_{ff} = \sqrt{\frac{3\pi}{16G\rho}}$$

c.f. Larson, R. D. 1985, MNRAS, 214, 379

Molecular gas mass needs to be concentrated to sheets or filamentary structures

My Proposals



Liu, H. B. et al. 2012, ApJ, 745, 61



A Tomographic View in the ¹³CO Velocity Channel Map



A Tomographic View in the ¹³CO Velocity Channel Map



OB Cluster-forming Regions in Actual Observations



W49A



Galvan-Madrid, R.. et al. 2013, A&A, 779, 121

G10.6-0.4



Liu, H. B. et al. 2012, ApJ, 745, 61

G10.6-0.4



G10.6-0.4





Flattened, gravitationally unstable rotating structure



Flattened, gravitationally unstable rotating structure



Flattened, gravitationally unstable rotating structure



Liu, H. B. et al. 2019, ApJ, 871, 185

Summary – What we do not understand Kennicutt-Schmidt Law & Gao-Solomon Relation Stellar Initial Mass Function Energetic and Kinematics in the Starforming Molecular Clouds

1. Role of turbulence

2. Role of magnetic field

3. Role of Feedback and cloud-cloud collision, galactic dynamics, etc.

The origin of SMBH and the formation of the M- σ relation